# ENVE2061 Basic Fluid Mechanics

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#### **FLUID STATICS**

**Pressure at a Point** 

**Basic Equation for Pressure Field** 

**Pressure Variation in a Fluid at Rest** 

Compressible Fluid Incompressible Fluid

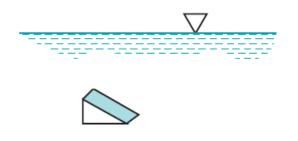
**Pressure Head** 

**Absolute & Gage Pressure** 

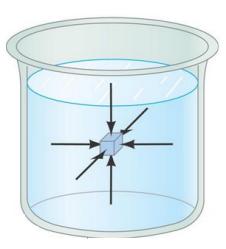
#### **FLUID STATICS**

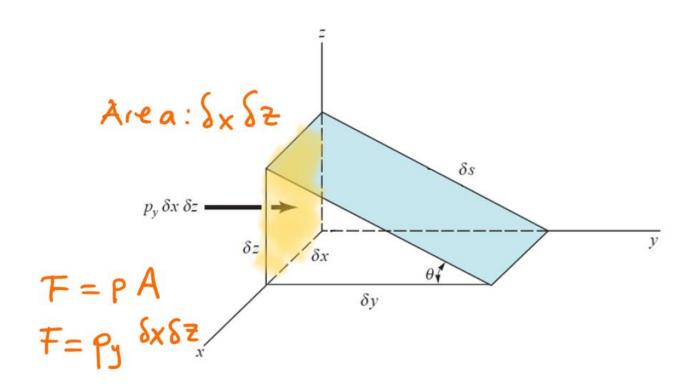
- determine the pressure at various locations in a fluid at rest.
- explain the concept of manometers and apply appropriate equations to determine pressures.
- calculate the hydrostatic pressure force on a plane or curved submerged surface.
- calculate the buoyant force and discuss the stability of floating or submerged objects.

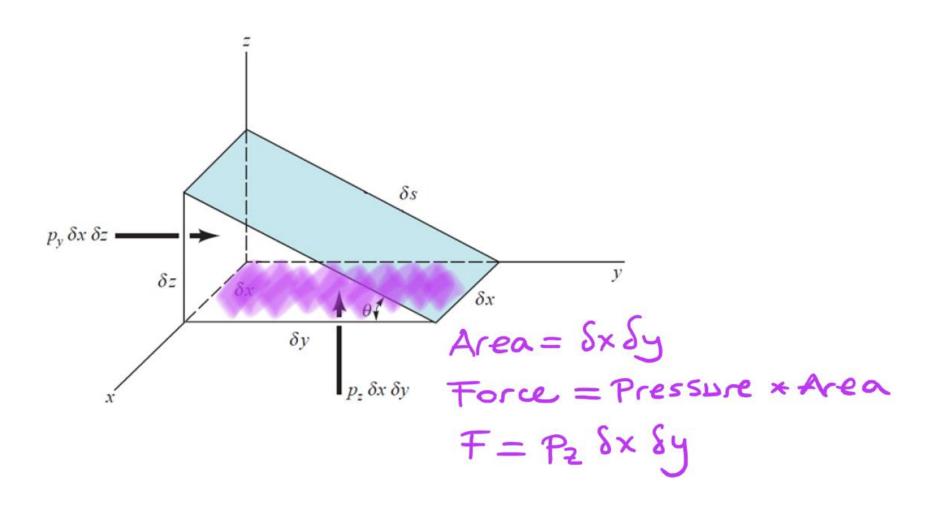
How the pressure at a point varies with the orientation of the plane passing through the point.

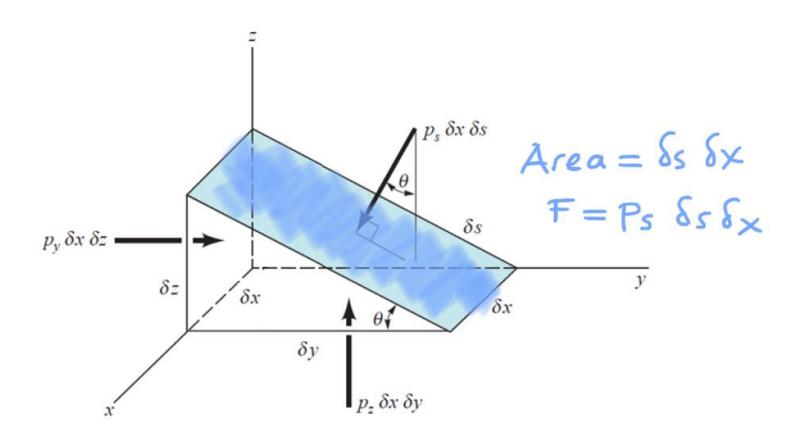


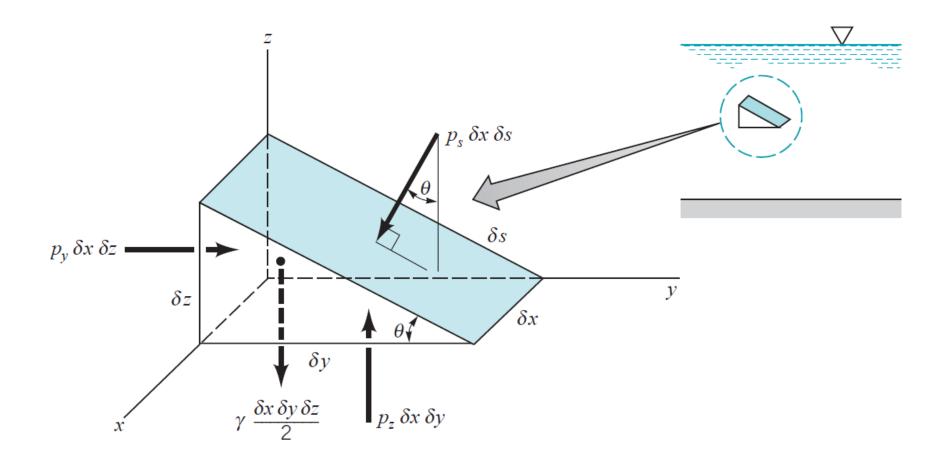
To answer this question, we consider the free-body diagram, that was obtained by removing a small triangular wedge of fluid from some arbitrary location within a fluid mass.







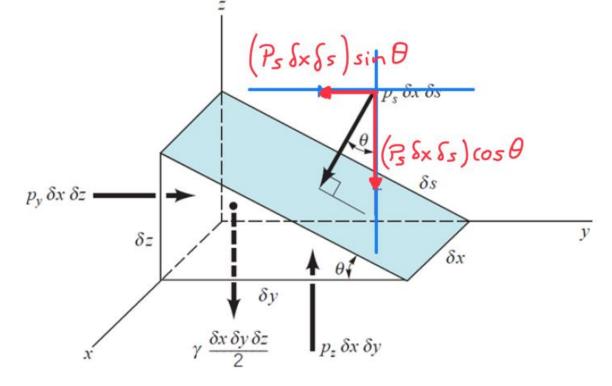




The equations of motion (Newton's second law, F = ma) in the y and z directions:

$$\sum F_y = p_y \, \delta x \, \delta z - \underline{p_s} \, \delta x \, \delta s \sin \theta$$
$$= \rho \, \frac{\delta x \, \delta y \, \delta z}{2} \, a_y$$

$$\sum F_z = p_z \, \delta x \, \delta y - \underline{p_s \, \delta x \, \delta s \cos \theta} - \gamma \, \frac{\delta x \, \delta y \, \delta z}{2}$$
$$= \rho \, \frac{\delta x \, \delta y \, \delta z}{2} \, a_z$$

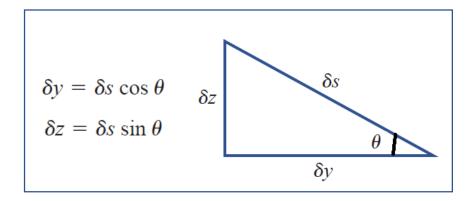


Average pressures on the faces:  $p_{s}$ ,  $p_{y}$ , and  $p_{z}$ 

Fluid specific weight:  $\gamma$ 

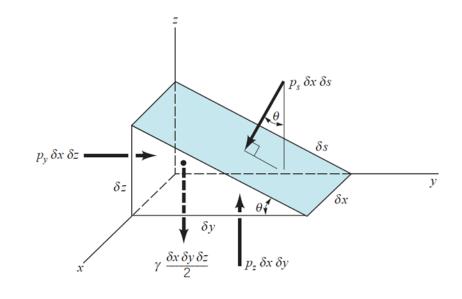
Fluid density:  $\rho$ 

Accelerations:  $a_y$ ,  $a_z$ 



$$\sum F_y = p_y \, \delta x \, \delta z - p_s \, \delta x \, \delta s \sin \theta$$
$$= \rho \, \frac{\delta x \, \delta y \, \delta z}{2} \, a_y$$

$$\sum F_z = p_z \, \delta x \, \delta y - p_s \, \delta x \, \delta s \cos \theta - \gamma \, \frac{\delta x \, \delta y \, \delta z}{2}$$
$$= \rho \, \frac{\delta x \, \delta y \, \delta z}{2} \, a_z$$



$$p_y - p_s = \rho a_y \frac{\delta y}{2}$$

$$p_z - p_s = (\rho a_z + \gamma) \frac{\delta z}{2}$$

Since we are really interested in what is happening at a point, we take the limit as  $\delta_x$ ,  $\delta_y$ , and  $\delta_z$  approach zero

$$p_y = p_s$$
  $p_z = p_s$ 

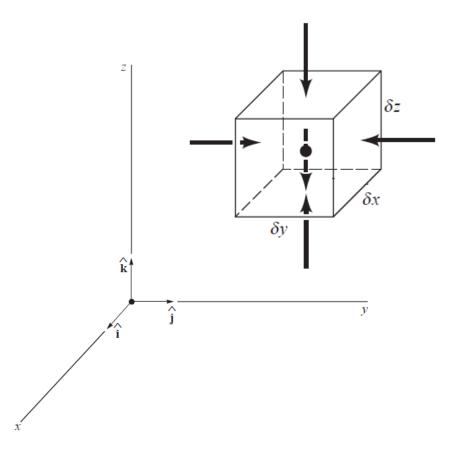
$$p_s = p_y = p_z$$

The pressure at a point in a fluid at rest, or in motion, is independent of direction as long as there are no shearing stresses present.

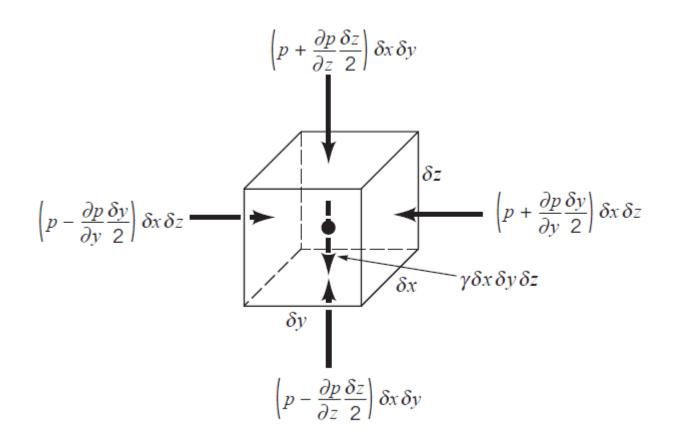
Pascal's law



How does the <u>pressure</u> in a fluid in which there are no shearing stresses <u>vary from point to point?</u>



There are two types of forces acting on this element: *surface forces* due to the pressure, and a *body force* equal to the weight of the element.

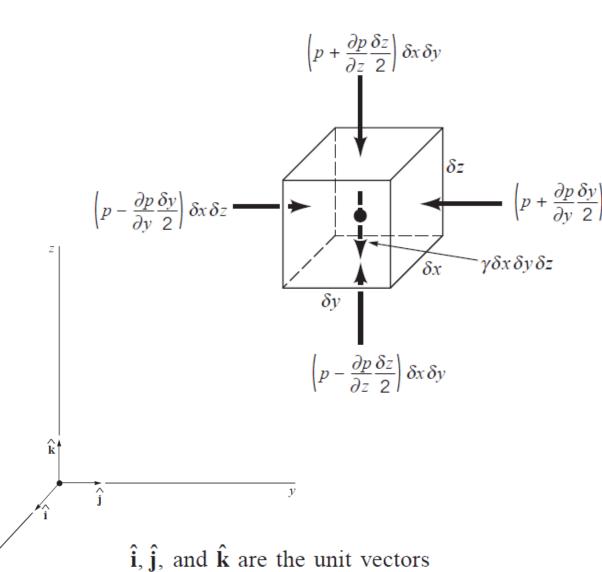


There are two types of forces acting on this element: *surface forces* due to the pressure, and a *body force* equal to the weight of the element.

The resultant surface force in the *y* direction

$$\delta F_{y} = \left(p - \frac{\partial p}{\partial y} \frac{\delta y}{2}\right) \delta x \, \delta z - \left(p + \frac{\partial p}{\partial y} \frac{\delta y}{2}\right) \delta x \, \delta z$$

$$\delta F_y = -\frac{\partial p}{\partial y} \, \delta x \, \delta y \, \delta z$$



along the coordinate axes

resultant surface forces

$$\delta F_y = -\frac{\partial p}{\partial y} \, \delta x \, \delta y \, \delta z$$

$$\delta F_x = -\frac{\partial p}{\partial x} \, \delta x \, \delta y \, \delta z$$

$$\delta F_z = -\frac{\partial p}{\partial z} \, \delta x \, \delta y \, \delta z$$

The resultant surface force acting on the element can be expressed in vector form as

$$\delta \mathbf{F}_{s} = \delta F_{x} \hat{\mathbf{i}} + \delta F_{y} \hat{\mathbf{j}} + \delta F_{z} \hat{\mathbf{k}}$$

$$\delta \mathbf{F}_{s} = \delta F_{x} \hat{\mathbf{i}} + \delta F_{y} \hat{\mathbf{j}} + \delta F_{z} \hat{\mathbf{k}}$$

$$\delta F_y = -\frac{\partial p}{\partial y} \, \delta x \, \delta y \, \delta z$$

$$\delta F_x = -\frac{\partial p}{\partial x} \, \delta x \, \delta y \, \delta z$$

$$\delta F_z = -\frac{\partial p}{\partial z} \, \delta x \, \delta y \, \delta z$$

$$\delta \mathbf{F}_{s} = -\left(\frac{\partial p}{\partial x}\,\hat{\mathbf{i}} + \frac{\partial p}{\partial y}\,\hat{\mathbf{j}} + \frac{\partial p}{\partial z}\,\hat{\mathbf{k}}\right) \delta x \,\delta y \,\delta z$$

$$\frac{\delta \mathbf{F}_s}{\delta x \, \delta y \, \delta z} = -\nabla p$$

resultant surface force per unit volume

$$\nabla(\ ) = \frac{\partial(\ )}{\partial x}\,\hat{\mathbf{i}} + \frac{\partial(\ )}{\partial y}\,\hat{\mathbf{j}} + \frac{\partial(\ )}{\partial z}\,\hat{\mathbf{k}}$$

symbol  $\nabla$  is the *gradient* or "del" vector operator

$$\frac{\partial p}{\partial x}\,\hat{\mathbf{i}} + \frac{\partial p}{\partial y}\,\hat{\mathbf{j}} + \frac{\partial p}{\partial z}\,\hat{\mathbf{k}} = \nabla p$$

$$\frac{\delta \mathbf{F}_s}{\delta x \, \delta v \, \delta z} = -\nabla p$$

resultant surface force per unit volume

$$-\delta \mathcal{W}\hat{\mathbf{k}} = -\gamma \, \delta x \, \delta y \, \delta z \, \hat{\mathbf{k}}$$

z axis is vertical, the weight of the element

Newton's second law, applied to the fluid element

$$\sum \delta \mathbf{F} = \delta m \mathbf{a}$$

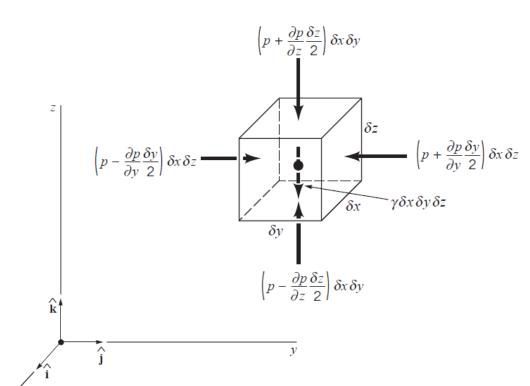
 $\Sigma$   $\delta F$  represents the resultant force acting on the element

$$\sum \delta \mathbf{F} = \delta \mathbf{F}_s - \delta^{\circ} W \hat{\mathbf{k}}$$
$$= -\nabla p \, \delta x \, \delta y \, \delta z - \gamma \, \delta x \, \delta y \, \delta z \, \hat{\mathbf{k}}$$

**a** is the acceleration of the element  $\delta m$  is the element mass,  $\rho \delta x \delta y \delta z$ 

$$\delta m \mathbf{a} = \rho \delta x \delta y \delta z \mathbf{a}$$

## **Basic Equation for Pressure Field**



$$-\nabla p - \gamma \hat{\mathbf{k}} = \rho \mathbf{a}$$

$$-\nabla p - \gamma \hat{\mathbf{k}} = \rho \mathbf{a}$$

general equation of motion for a fluid in which there are no shearing stresses

For a fluid at rest  $\mathbf{a} = 0$ 

$$\nabla p + \gamma \hat{\mathbf{k}} = 0 \qquad \nabla p = -\gamma \hat{\mathbf{k}}$$

$$\frac{\partial p}{\partial x}\,\hat{\mathbf{i}} + \frac{\partial p}{\partial y}\,\hat{\mathbf{j}} + \frac{\partial p}{\partial z}\,\hat{\mathbf{k}} = \nabla p = -\gamma\hat{\mathbf{k}}$$

$$\frac{\partial p}{\partial x} = 0$$
  $\frac{\partial p}{\partial y} = 0$   $\frac{\partial p}{\partial z} = -\gamma$ 

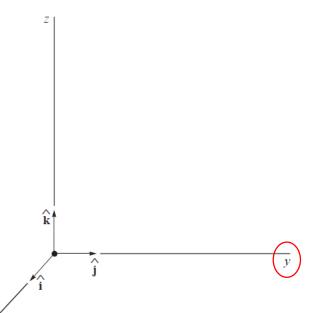
pressure does not depend on *x* or *y* 

$$\nabla(\ ) = \frac{\partial(\ )}{\partial x}\,\hat{\mathbf{i}} + \frac{\partial(\ )}{\partial y}\,\hat{\mathbf{j}} + \frac{\partial(\ )}{\partial z}\,\hat{\mathbf{k}}$$

symbol  $\nabla$  is the *gradient* or "del" vector operator

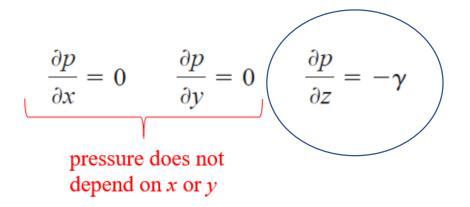
$$\frac{\partial p}{\partial x}\,\hat{\mathbf{i}} + \frac{\partial p}{\partial y}\,\hat{\mathbf{j}} + \frac{\partial p}{\partial z}\,\hat{\mathbf{k}} = \nabla p$$

## Pressure Variation in a Fluid at Rest



as we move from point to point in a horizontal plane (any plane parallel to the x–y plane), the pressure does not change

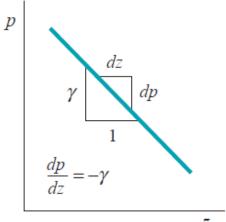
#### **Pressure Variation in a Fluid at Rest**



Since *p* depends only on *z*, equation can be written as the ordinary differential equation

$$\frac{dp}{dz} = -\gamma$$

pressure decreases as we move upward in a fluid at rest



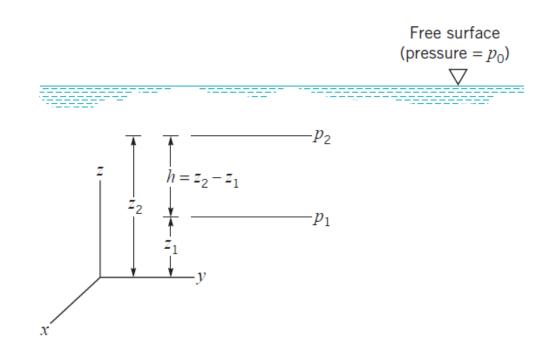
### Pressure Variation in a Fluid at Rest – Incompressible Fluid

$$\frac{dp}{dz} = -\gamma$$

$$\int_{p_1}^{p_2} dp = -\gamma \int_{z_1}^{z_2} dz$$

$$p_1 - p_2 = \gamma h$$

$$p_1 = \gamma h + p_2$$



This type of pressure distribution is commonly called **hydrostatic distribution**.

in an incompressible fluid at rest the pressure varies linearly with depth

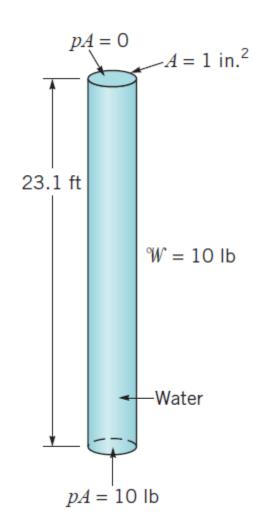
### Pressure Variation in a Fluid at Rest - Incompressible Fluid

#### PRESSURE HEAD

$$p_1 = \gamma h + p_2$$

$$h = \frac{p_1 - p_2}{\gamma}$$

h is called the *pressure head* and is interpreted as the height of a column of fluid of specific weight  $\gamma$  required to give a pressure difference  $p_1 - p_2$ .



Pressure difference 10 psi Express as pressure head:

pressure head as 23.1 ft of water  $(\gamma = 62.4 \text{ lb/ft}^3)$ 

518 mm of Hg 
$$(\gamma = 133 \text{ kN/m}^3)$$

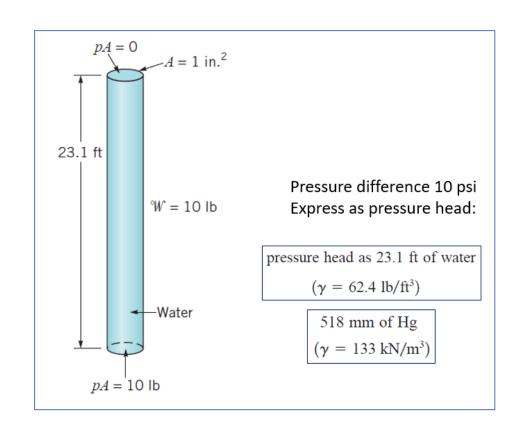
## Pressure Variation in a Fluid at Rest – Incompressible Fluid PRESSURE HEAD

### Pressure Variation in a Fluid at Rest – Incompressible Fluid

#### PRESSURE HEAD

$$h_{120} = \frac{\Delta P}{4 + 20} = \frac{10 + 16}{162 \cdot 4} = \frac{12 i \pi}{162 \cdot 4} = 23.1 \text{ ft}$$

$$= \frac{10 \times 144}{62 \cdot 4} = 23.1 \text{ ft}$$



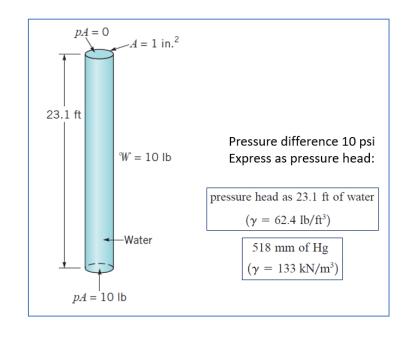
### Pressure Variation in a Fluid at Rest – Incompressible Fluid

#### PRESSURE HEAD

$$h_{y} = \frac{\Delta P}{M_{y}} = \frac{10 \frac{16}{112} \left[\frac{12i\pi}{1+12}\right]^{2} \frac{4.4482 \text{ M}}{1 \text{ M}} \cdot \frac{16\pi}{1000 \text{ M}} \left[\frac{144}{0.3048^{2}}\right]}{(0*12^{2}*4.4482* + \frac{1}{1000}*\frac{1}{0.3048^{2}})} = 0.518 \text{ m}$$

$$= \frac{133}{133} = 0.518 \text{ m}$$

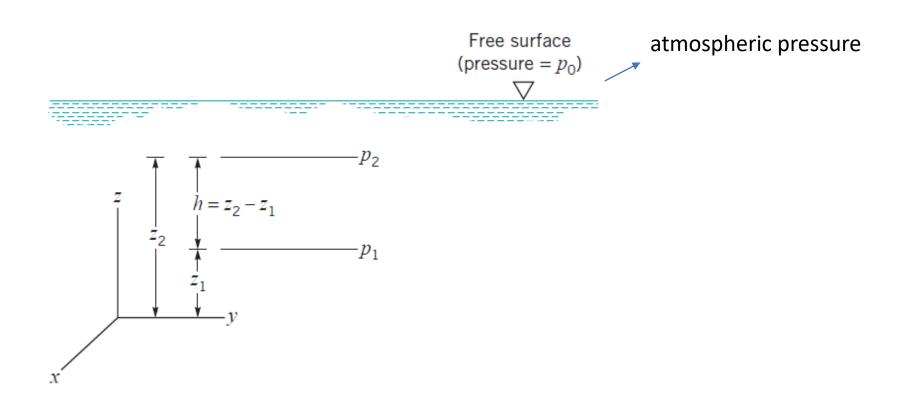
$$= \frac{133}{133} = 0.518 \text{ m}$$



$$0.518 \text{ m} = 518 \text{ m}$$

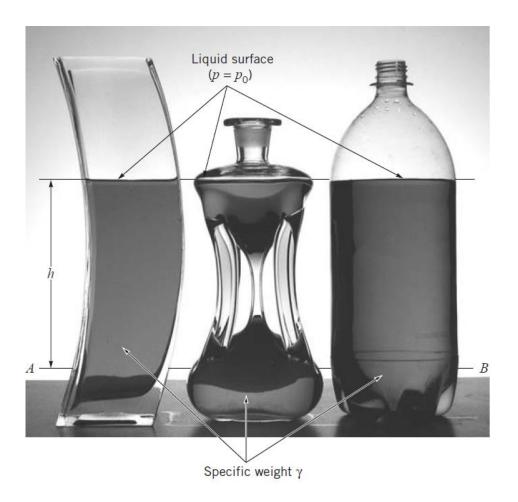
$$12 \text{ in} = 1 \text{ ft}$$
 $1 \text{ ft} = 0.3048 \text{ m}$ 
 $1.16 = 4.4482 \text{ N}$ 

# FLUID STATICS - Pressure Variation in a Fluid at Rest Incompressible Fluid



## FLUID STATICS - Pressure Variation in a Fluid at Rest

## Incompressible Fluid





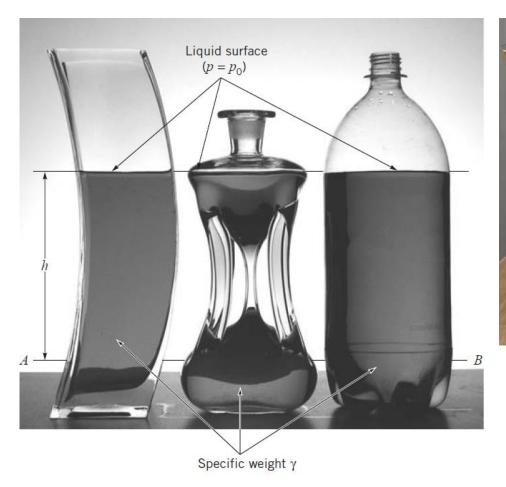
$$p = \gamma h + p_0$$





## FLUID STATICS - Pressure Variation in a Fluid at Rest

## Incompressible Fluid





$$p = \gamma h + p_0$$

## FLUID STATICS - Pressure Variation in a Fluid at Rest Incompressible Fluid

#### EXAMPLE 2.1 Pressure-Depth Relationship

**GIVEN** Because of a leak in a buried gasoline storage tank, water has seeped in to the depth shown in Fig. E2.1. The specific gravity of the gasoline is SG = 0.68.

FIND Determine the pressure at the gasoline-water interface and at the bottom of the tank. Express the pressure in units of 1b/ft², 1b/in.², and as a pressure head in feet of water.

#### SOLUTION

Since we are dealing with liquids at rest, the pressure distribution will be hydrostatic, and therefore the pressure variation can be found from the equation:

$$p = \gamma h + p_0$$

With  $p_0$  corresponding to the pressure at the free surface of the gasoline, then the pressure at the interface is

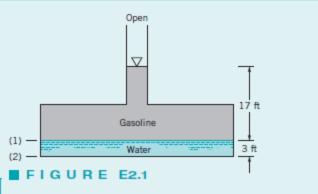
$$p_1 = SG\gamma_{\text{H}_2\text{O}}h + p_0$$
  
= (0.68)(62.4 lb/ft<sup>3</sup>)(17 ft) + p\_0  
= 721 + p\_0 (lb/ft<sup>2</sup>)

If we measure the pressure relative to atmospheric pressure (gage pressure), it follows that  $p_0=0$ , and therefore

$$p_1 = 721 \text{ lb/ft}^2$$
 (Ans)

$$p_1 = \frac{721 \text{ lb/ft}^2}{144 \text{ in }^2/\text{ft}^2} = 5.01 \text{ lb/in.}^2$$
 (Ans)

$$\frac{p_1}{2m_0} = \frac{721 \text{ lb/ft}^2}{62.4 \text{ lb/ft}^3} = 11.6 \text{ ft}$$
 (Ans)



It is noted that a rectangular column of water 11.6 ft tall and 1 ft<sup>2</sup> in cross section weighs 721 lb. A similar column with a 1-in.<sup>2</sup> cross section weighs 5.01 lb.

We can now apply the same relationship to determine the pressure at the tank bottom; that is,

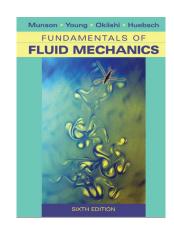
$$p_2 = \gamma_{\text{H}_2\text{O}} h_{\text{H}_2\text{O}} + p_1$$
  
=  $(62.4 \text{ lb/ft}^3)(3 \text{ ft}) + 721 \text{ lb/ft}^2$  (Ans)  
=  $908 \text{ lb/ft}^2$ 

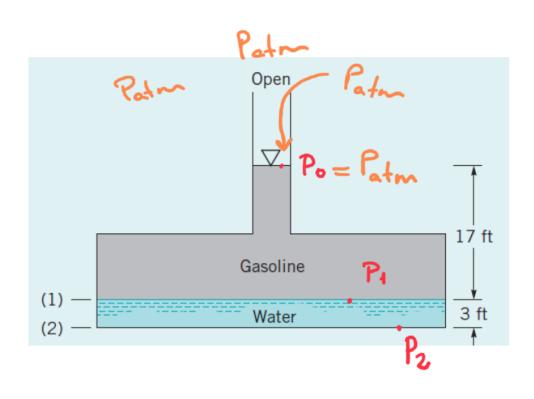
$$p_2 = \frac{908 \text{ lb/ft}^2}{144 \text{ in.}^2/\text{ft}^2} = 6.31 \text{ lb/in.}^2$$
 (Ans)

$$\frac{p_2}{\gamma_{\text{H,O}}} = \frac{908 \text{ lb/ft}^2}{62.4 \text{ lb/ft}^3} = 14.6 \text{ ft}$$
 (Ans)

(Ans)

COMMENT Observe that if we wish to express these pressures in terms of absolute pressure, we would have to add the local atmospheric pressure (in appropriate units) to the previous results. A further discussion of gage and absolute pressure is given in Section 2.5.





$$P_{1} = P_{0} + \frac{1}{3} h_{9}$$

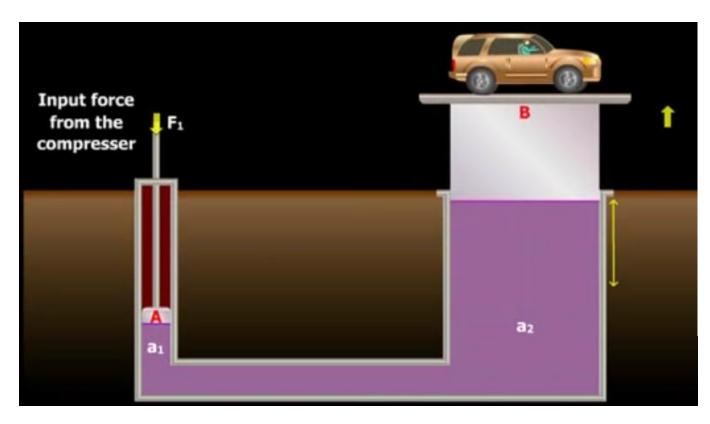
$$P_{2} = P_{1} + \frac{1}{3} h_{2} 0 h_{12} 0$$

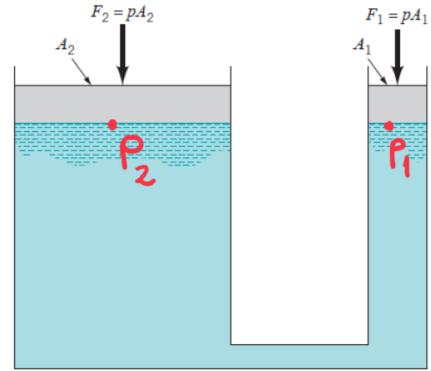
$$P_{2} = P_{0} + \frac{1}{3} h_{9} + \frac{1}{3} h_{2} 0 h_{12} 0$$

$$P_{3} = P_{0} + \frac{1}{3} h_{9} + \frac{1}{3} h_{12} 0 h_{12} 0$$
ine  $P_{4}$ 

## FLUID STATICS - Pressure Variation in a Fluid at Rest

Incompressible Fluid
Transmission of Fluid Pressure





p acting on the faces of both pistons is the same

$$P_1 = P_2 = P$$

$$F_2 = (A_2/A_1)F_1$$

### FLUID STATICS - Pressure Variation in a Fluid at Rest

## Incompressible Fluid Transmission of Fluid Pressure

#### Example 2.1

The diameters of cylindrical pistons A and B are 3 cm and 20 cm, respectively. The faces of the pistons are at the same elevation, and the intervening passages are filled with an incompressible hydraulic oil. A force P of 100 N is applied at the end of the lever, as shown in Figure 2.3. What weight W can the hydraulic jack support?

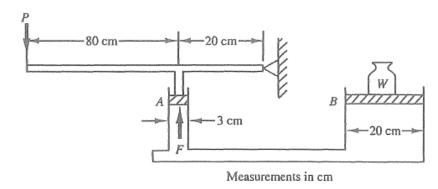


Figure 2.3 Hydraulic jack

#### Solution

Balancing the moments produced by P and F around the pin connection yields

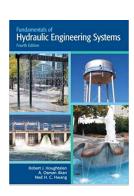
$$(100 \text{ N})(100 \text{ cm}) = F(20 \text{ cm})$$

Thus,

$$F = 500 \, \text{N}$$

From Pascal's law, the pressure  $P_A$  applied at A is the same as that of  $P_B$  applied at B. Therefore,

$$P_A = \frac{F}{[(\pi \cdot 3^2)/4] \text{ cm}^2} \qquad P_B = \frac{W}{[(\pi \cdot 20^2)/4] \text{ cm}^2}$$
$$\frac{500 \text{ N}}{7.07 \text{ cm}^2} = \frac{W}{314 \text{ cm}^2}$$
$$\therefore W = 500 \text{ N} \left(\frac{314 \text{ cm}^2}{7.07 \text{ cm}^2}\right) = 2.22 \times 10^4 \text{ N}$$



# FLUID STATICS - Pressure Variation in a Fluid at Rest Compressible Fluid

Gases such are compressible fluids since the density of the gas can change significantly with changes in pressure & temperature.

$$\frac{dp}{dz} = -\gamma$$

## FLUID STATICS - Pressure Variation in a Fluid at Rest

### Compressible Fluid

Ideal gas law 
$$\rho = \frac{p}{RT} \qquad (\gamma = \rho g)$$
 
$$\rho g = \frac{gp}{RT}$$
 
$$\rho g = \frac{gp}{RT}$$

$$-\gamma = -\frac{gp}{RT}$$

$$\frac{dp}{dz} = -\gamma$$

$$\frac{dp}{dz} = -\frac{gp}{RT}$$

# FLUID STATICS - Pressure Variation in a Fluid at Rest Compressible Fluid

$$\frac{dp}{dz} = -\frac{gp}{RT}$$

$$\int_{p_1}^{p_2} \frac{dp}{p} = \ln \frac{p_2}{p_1} = -\frac{g}{R} \int_{z_1}^{z_2} \frac{dz}{T}$$

$$p_2 = p_1 \exp\left[-\frac{g(z_2 - z_1)}{RT_0}\right]$$

## FLUID STATICS - Pressure Variation in a Fluid at Rest Compressible Fluid

#### EXAMPLE 2.2 Incompressible and Isothermal Pressure-Depth Variations

GIVEN In 2007 the Burj Dubai skyscraper being built in the United Arab Emirates reached the stage in its construction where it became the world's tallest building. When completed it is expected to be at least 2275 ft tall, although its final height remains a secret

**FIND** (a) Estimate the ratio of the pressure at the projected 2275-ft top of the building to the pressure at its base, assuming the air to be at a common temperature of 59 °F. (b) Compare the pressure calculated in part (a) with that obtained by assuming the air to be incompressible with  $\gamma = 0.0765 \text{ lb/ft}^3$  at 14.7 psi (abs) (values for air at standard sea level conditions).

#### SOLUTION

For the assumed isothermal conditions, and treating air as a compressible fluid, Eq. 2.10 can be applied to yield

$$\frac{p_2}{p_1} = \exp\left[-\frac{g(z_2 - z_1)}{RT_0}\right] 
= \exp\left\{-\frac{(32.2 \text{ ft/s}^2)(2275 \text{ ft})}{(1716 \text{ ft} \cdot \text{lb/slug} \cdot ^\circ\text{R})[(59 + 460)^\circ\text{R}]}\right\} 
= 0.921$$
(Ans)

If the air is treated as an incompressible fluid we can apply Eq. 2.5. In this case

$$p_2=p_1-\gamma(z_2-z_1)$$

Of

$$\frac{p_2}{p_1} = 1 - \frac{\gamma(z_2 - z_1)}{p_1}$$

$$= 1 - \frac{(0.0765 \text{ lb/}\Re^3)(2275 \text{ ft})}{(14.7 \text{ lb/in.}^2)(144 \text{ in.}^2/\Re^2)} = 0.918 \quad \text{(Ans)}$$

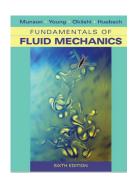
COMMENTS Note that there is little difference between the two results. Since the pressure difference between the bottom and top of the building is small, it follows that the variation in fluid density is small and, therefore, the compressible



■ FIGURE E2.2 (Figure courtesy of Emaar Properties, Dubai, UAE.)

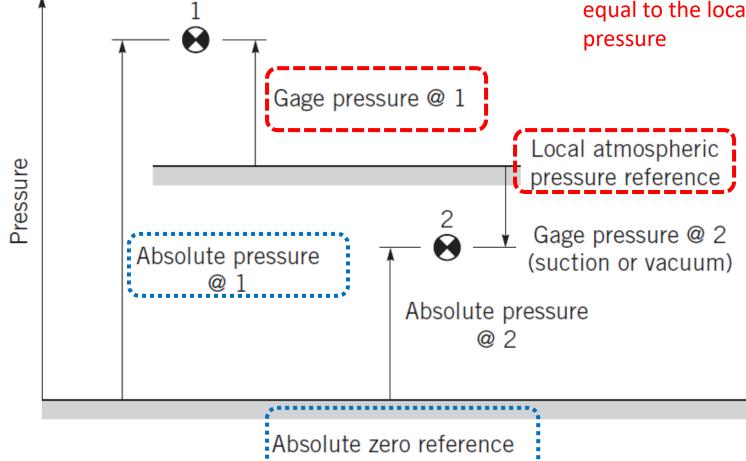
fluid and incompressible fluid analyses yield essentially the same result.

We see that for both calculations the pressure decreases by approximately 8% as we go from ground level to the top of this tallest building. It does not require a very large pressure difference to support a 2275-ft-tall column of fluid as light as air. This result supports the earlier statement that the changes in pressures in air and other gases due to elevation changes are very small, even for distances of hundreds of feet. Thus, the pressure differences between the top and bottom of a horizontal pipe carrying a gas, or in a gas storage tank, are negligible since the distances involved are very small.

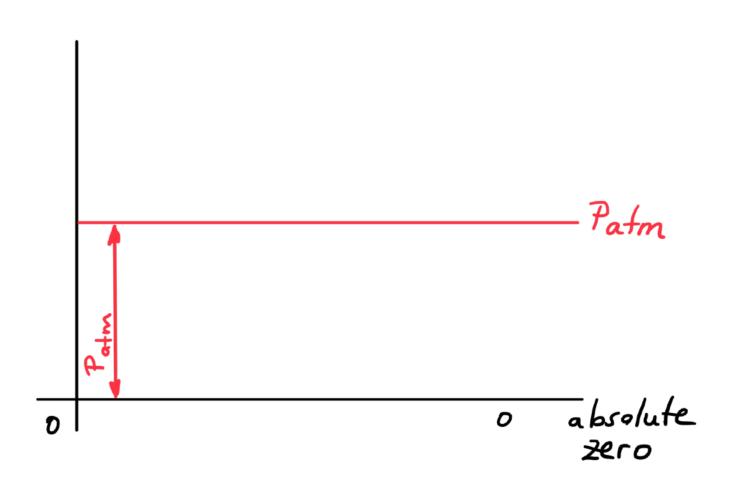


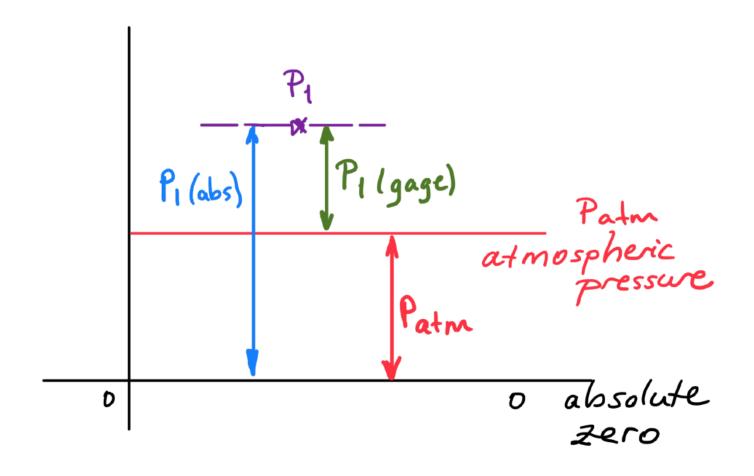
# FLUID STATICS — Measurement of Pressure Absolute Pressure & Gage Pressure

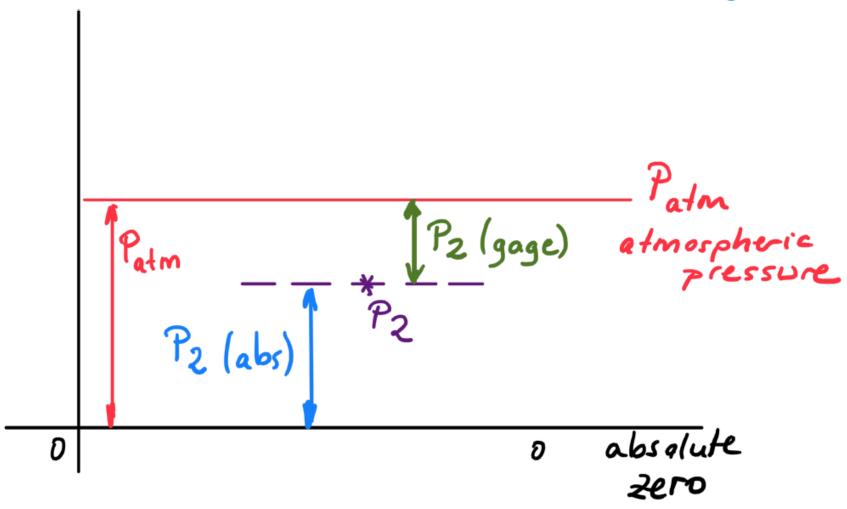
Gage pressure is measured relative to the local atmospheric pressure. Thus a gage pressure of zero corresponds to a pressure that is equal to the local atmosperic pressure

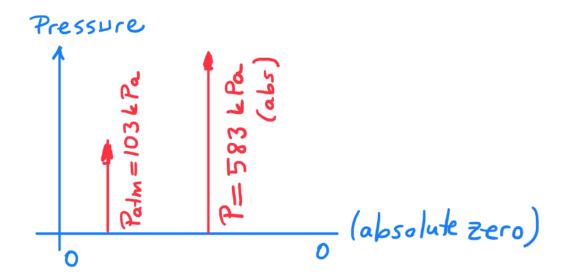


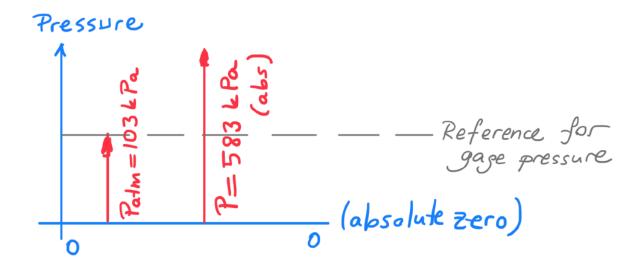
Absolute pressure is measured relative to a perfect vacuum (absolute zero pressure)





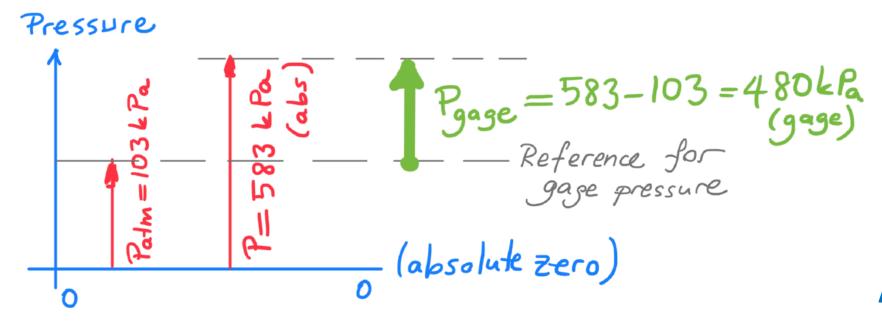


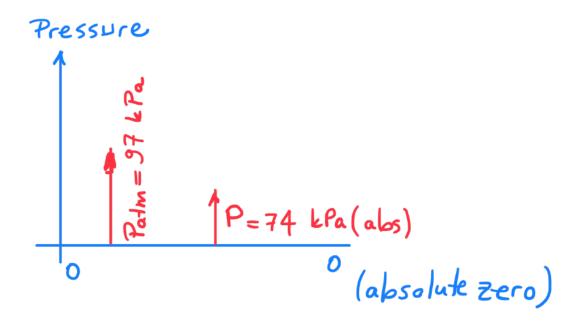




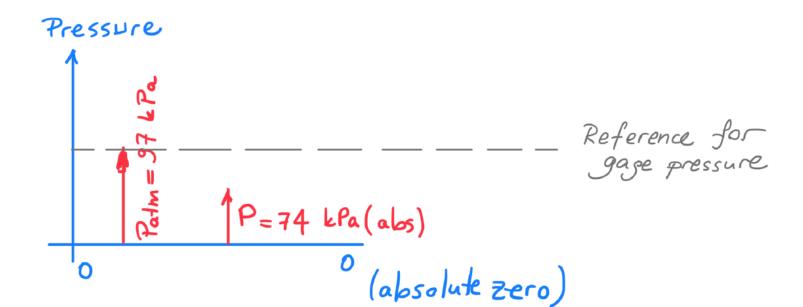
	Given	Express Rest		
Pressure		$p_{\rm atm}$		As:
583	kPa(abs)	103	kPa(abs)	Gage pressure

	Given	Express R		Express Result
Pressure		$p_{\rm atm}$		As:
583	kPa(abs)	103	kPa(abs)	Gage pressure





Given Pressure		$p_{ m atm}$		Express Result As:
74	kPa(abs)	97	kPa(abs)	Gage pressure



_	iven ssure	$p_{ m atm}$		Express Result As:
74	kPa(abs)	97	kPa(abs)	Gage pressure

